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# Correlation functions for a periodic box-ball system 

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#### Abstract

We investigate correlation functions in a periodic box-ball system. For the two-point functions of short distance, we give explicit formulae obtained by combinatorial methods. We give expressions for general N -point functions in terms of ultradiscrete theta functions.


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## 1. Introduction

Quantum integrable systems such as quantum integrable spin chains and solvable lattice models are systems whose Hamiltonians or transfer matrices can be diagonalized and for which eigenstates or free energies can be explicitly obtained [1]. However, to investigate physical properties of these systems, such as e.g. the linear response to external forces, we further need to evaluate correlation functions for these systems. This is one of the main problems in the field of quantum integrable systems and, in fact, obtaining correlation functions is fairly difficult even for the celebrated $X X Z$ model or the six-vertex model [2].

A periodic box-ball system (PBBS) is a soliton cellular automaton obtained by ultradiscretizing the KdV equation [3, 4]. It can also be obtained at the $q \rightarrow 0$ limit of the generalized six-vertex model [5, 6]. Hence, from the viewpoint of quantum integrable lattice models, it is interesting and may actually give some new insights into the correlation functions of the vertex models themselves to obtain correlation functions of the PBBS. In this paper, we give expressions for $N$-point functions for the PBBS, using combinatorial methods and the solution for the PBBS expressed in terms of the ultradiscrete theta functions.

The PBBS can be defined as follows. Let $L \geqslant 3$ and let $\Omega_{L}=\{f \mid f:[L] \rightarrow$ $\{0,1\}$ such that $\left.\sharp f^{-1}(\{1\})<L / 2\right\}$, where $[L]=\{1,2, \ldots, L\}$. When $f \in \Omega_{L}$ is represented as a sequence of 0 s and 1 s , we write

$$
f(1) f(2) \cdots f(L)
$$

The mapping $T_{L}: \Omega_{L} \rightarrow \Omega_{L}$ is defined as follows (see figure 1).


Figure 1. Definition of $T_{L}$ for $f \in \Omega_{L}$.

1. In the sequence $f$, find a pair of positions $n$ and $n+1$ such that $f(n)=1$ and $f(n+1)=0$, and mark them; repeat the same procedure until all such pairs are marked. Note that we always use the convention that the position $n$ is defined in $[L]$, i.e. $n+L \equiv n$.
2. Skipping the marked positions we get a subsequence of $f$; for this subsequence, repeat the same process of marking positions, so that we get another marked subsequence.
3. Repeat step 2 until a subsequence consisting only of 0 s is obtained. A typical situation is depicted in figure 1. After these preparatory processes, change all values at the marked positions simultaneously. One thus obtains the sequence $T_{L} f$.
Sometimes we shall write $T_{L}^{t} f$ for $\underbrace{T_{L}\left(\cdots\left(T_{L}\left(T_{L}\right.\right.\right.}_{t} f)))$. The pair $\left(\Omega_{L}, T_{L}\right)$ is called a PBBS of length $L[4,7]$. An element of $\Omega_{L}$ is called a state, and the mapping $T_{L}$ the time evolution.

The conserved quantities of the PBBS are defined as follows. Let $Q_{j}(t)$ be the number 10 pairs in $T_{L}^{t} f$ marked at the $j$ th step in the definition of the mapping $T_{L}$. Then we obtain a nonincreasing sequence of positive integers, $Q_{j}(t)(j=1,2, \ldots, m)$. This sequence is conserved in time, that is,

$$
Q_{j}(t)=Q_{j}(t+1) \equiv Q_{j} \quad(j=1,2, \ldots, m)
$$

For example, for $f$ in figure $1,\left(Q_{1}, Q_{2}, Q_{3}, Q_{4}\right)=(3,3,2,1)$. As the sequence ( $Q_{1}, Q_{2}, \ldots, Q_{m}$ ) is nonincreasing, we can associate a Young's diagram with it by considering $Q_{j}$ as the number of squares in the $j$ th column of the diagram. The lengths of the rows are also weakly decreasing positive integers. Let the distinct row lengths be $P_{1}>P_{2}>\cdots>P_{\ell}$ and let $n_{j}$ be the number of times the length $P_{j}$ appears. The set $\left\{P_{j}, n_{j}\right\}_{j=1}^{\ell}$ is another expression for the conserved quantities of the PBBS.

An $N$-point function of the PBBS with $M$ balls may be defined as follows.

$$
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle:=\frac{1}{Z_{H}} \sum_{f \in \Omega_{L ; M}} \mathrm{e}^{\sum_{k=1}^{L} H_{k}(f)} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right)
$$

where $\Omega_{L ; M}:=\left\{f \in \Omega_{L} \mid \sharp f^{-1}(\{1\})=M\right\}, Z_{H}:=\sum_{f \in \Omega_{L ; M}} \mathrm{e}^{\sum_{k=1}^{L} H_{k}(f)}$ and $H_{k}(f)$ is the $k$ th energy of the state $f$, which is proportional to the number of $k$ th arc lines defined when determining the time evolution rule [4], or the $k$ th conserved quantity of the PBBS [8]. (Note that $H_{k}(f)$ is essentially equal to the energy function for the transfer matrix of the crystal lattice models with $k+1$ states on a vertical link [5, 7].) Noticing the fact that $\Omega_{L ; M}=\bigsqcup_{Y} \Omega_{Y}$,

$$
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle=\frac{1}{Z_{H}} \sum_{Y} \sum_{f \in \Omega_{Y}} \mathrm{e}^{\sum_{k=1}^{L} H_{k}(f)} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right),
$$

where $Y$ are partitions of $M$ corresponding to the conserved quantities of the PBBS (see section 2). Since, for $f_{i} \in \Omega_{Y_{i}}(i=1,2), \forall k, H_{k}\left(f_{1}\right)=H_{k}\left(f_{2}\right)(k=1,2,3, \ldots)$ implies
$Y_{1}=Y_{2}$, and vice versa, by choosing a state $f_{Y}$ in $\Omega_{Y}$, we can write

$$
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle=\frac{1}{Z_{H}} \sum_{Y} \mathrm{e}^{\sum_{k=1}^{L} H_{k}\left(f_{Y}\right)} \sum_{f \in \Omega_{Y}} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right)
$$

Thus, to obtain correlation functions of PBBS, we only have to evaluate those on the set $\Omega_{Y}$ :

$$
\begin{equation*}
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle_{Y}:=\frac{1}{\left|\Omega_{Y}\right|} \sum_{f \in \Omega_{Y}} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right) \tag{1}
\end{equation*}
$$

We also point out that if we put $\forall k, \forall f, H_{k}(f)=0, N$-point functions become trivial:

$$
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle=\frac{L_{-N} C_{M-N}}{{ }_{L} C_{M}}=\frac{M(M-1) \cdots(M-N+1)}{L(L-1) \cdots(L-N+1)} .
$$

In the following sections, we shall evaluate (1).
First we summarize some useful properties of the PBBS. We say that $f$ has (or that there is) a 10 -wall at position $n$ if $f(n-1)=1$ and $f(n)=0$. Let the number of the 10 -walls be $s$ and let their positions be denoted by $a_{1}>a_{2}>\cdots>a_{s}$. Then, we have the following proposition.

## Proposition 1 ([9]).

$$
\begin{align*}
& \left(T_{L}^{t} f\right)(n)=\eta_{n+1}^{t-1}-\eta_{n+1}^{t}-\eta_{n}^{t-1}+\eta_{n}^{t} \\
& \eta_{n}^{t}=\max _{\substack{m_{i} \in \mathbb{Z} \\
i \in[s]}}\left[\sum_{i=1}^{s} m_{i}\left(b_{i}+t W_{i}-n\right)-\sum_{i=1}^{\ell} \sum_{j=1}^{\ell} m_{i} \Xi_{i j} m_{j}\right]  \tag{2}\\
& b_{i}=a_{i}+\sum_{j=1}^{i-1} 2 \min \left\{W_{i}, W_{j}\right\}+W_{i}+\frac{Z_{i}}{2},  \tag{3}\\
& \Xi_{i j}=\frac{Z_{i}}{2} \delta_{i j}+\min \left\{W_{i}, W_{j}\right\} \\
& Z_{i}=L-\sum_{j=1}^{s} 2 \min \left\{W_{i}, W_{j}\right\}
\end{align*}
$$

where $W_{i}$ denotes the amplitude of the 'soliton' corresponding to $a_{i}$ obtained by the procedure explained in [9].

The set $\left\{W_{i}\right\}_{i=1}^{s}$ consists of quantities of the PBBS and $\eta_{n}^{t}$ is the ultradiscrete theta function [10]. We shall use proposition 1 for determining $N$-point functions in section 3.

Next we introduce two procedures which are important in this paper. For a given $f \in \Omega_{L}$, a state $E f=E(f)$ is defined to be

The mapping $E: \Omega_{L} \rightarrow \Omega_{L-2 s}$ is called the 10 -elimination. $E f$ is a subsequence of $f$ obtained by eliminating all the 10 -walls in $f$ simultaneously. For example,


```
Ef=001111 0001111 0000000111 }00
    =00111100011110000000111010001111011000111000000000.
```

Its inverse process is called the 10 -insertion, $I\left(j_{1}, j_{2}, \ldots, j_{d}\right)=I_{2} \circ I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right)$ : $\Omega_{L} \rightarrow \Omega_{L+2(d+s)}$, where $s$ is the number of 10 -walls in $f \in \Omega_{L}$. The 10 -insertion is defined as follows: shifting the origin if necessary, we can assume that $f(L)=0$. For $\left\{j_{1}, j_{2}, \ldots, j_{d}\right\}\left(1<j_{1}<j_{2}<\cdots<j_{d} \leqslant L+d\right)$, the mapping $I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right): \Omega_{L} \rightarrow$ $\Omega_{L+2 d}$ is defined as
$\left(I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right) f\right)(n)$

$$
= \begin{cases}1 & \left(n=L+2 d-j_{k}-k+1\right) \\ 0 & \left(n=L+2 d-j_{k}-k+2\right) \\ f(n) & \left(1 \leqslant n \leqslant L+d-j_{d}\right) \\ f(n-2(d-k+1)) & \left(L+2 d-j_{k}-k+3 \leqslant n \leqslant L+2 d-j_{k-1}-k+1\right) \\ f(n-2 d) & \left(L+2 d-j_{1}+2 \leqslant n \leqslant L+2 d\right)\end{cases}
$$

where $k \in[d]$; furthermore, $I_{2}: \Omega_{L+2 d} \rightarrow \Omega_{L+2(d+s)}$ is defined to be

$$
\left(I_{2} f^{\prime}\right)(n)= \begin{cases}1 & \left(n=g_{k}+2(s-k)+1\right) \\ 0 & \left(n=g_{k}+2(s-k)+2\right) \\ f^{\prime}(n) & \left(1 \leqslant n \leqslant g_{s}\right) \\ f^{\prime}(n-2(s-k+1)) & \left(g_{k}+2(s-k)+3 \leqslant n \leqslant g_{k-1}+2(s-k)+2\right) \\ f^{\prime}(n-2 s) & \left(g_{1}+2 s+1 \leqslant n \leqslant L+2(d+s)\right)\end{cases}
$$

where $k \in[s], f^{\prime} \equiv I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right) f \in \Omega_{L+2 d}$ and

$$
\begin{align*}
& g_{k}^{\prime}=\max \left\{m \in[L+d] \mid m=a_{k}-1+\sharp\left\{r \in[d] \mid L+d-j_{r}+1<m\right\}\right\} \\
& g_{k}=g_{k}^{\prime}+\sharp\left\{r \in[d] \mid L+d-j_{r}+1<g_{k}^{\prime}\right\} \tag{4}
\end{align*}
$$

For example,
$f=0011100111000001101000111000000$,

$$
\begin{align*}
I_{1}(3,11,25) f & =001110011 * 1000001101000 * 1110000 * 00  \tag{5}\\
& =001110011 \underline{1} 10000011010001 \underline{1} 1110000 \underline{1000} \tag{6}
\end{align*}
$$

$I(3,11,25) f=00111100011 \underline{101} 100000011 \boxed{10} 01000101111000001000$,
where $\underline{10}$ and 10 , respectively, denote the inserted 10 at $f \mapsto I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right) f$ and $I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right) f \mapsto I_{2}\left(I_{1}\left(j_{1}, j_{2}, \ldots, j_{d}\right) f\right)$.

In the above example, $g_{k}^{\prime}$ is the position of 1 found in the $k$ th 10 -wall in the sequence (5), and $g_{k}$ is the position of 1 found in the $k$ th 10 -wall in the sequence (6). That is, $\left\{g_{k}^{\prime}\right\}_{k=1}^{5}=\{27,20,18,11,5\}$ and $\left\{g_{k}\right\}_{k=1}^{5}=\{29,21,19,12,5\}$.

## 2. One- and two-point functions obtained by combinatorial methods

We assume that $Y$ denoting the conserved quantities of $f \in \Omega_{Y}$ is the partition

$$
(\underbrace{P_{1}, P_{1}, \ldots, P_{1}}_{n_{1}}, \underbrace{P_{2}, P_{2}, \ldots, P_{2}}_{n_{2}}, \ldots, \underbrace{P_{\ell}, P_{\ell}, \ldots, P_{\ell}}_{n_{\ell}}),
$$

where $P_{1}>P_{2}>\cdots>P_{\ell} \geqslant 1$. Note that $Y$ is a partition of $M$, i.e. $M=\sum_{i=1}^{\ell} n_{i} P_{i}$. As mentioned in section 1, we consider $N$-point functions (1) of the PBBS:

$$
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle_{Y}=\frac{1}{\left|\Omega_{Y}\right|} \sum_{f \in \Omega_{Y}} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right)
$$

The value of $\left|\Omega_{Y}\right|$ is already known.
Proposition 2 ([11]).

$$
\left|\Omega_{Y}\right|=\frac{L}{L_{0}}\binom{L_{0}+n_{1}-1}{n_{1}}\binom{L_{1}+n_{2}-1}{n_{2}} \ldots\binom{L_{\ell-1}+n_{\ell}-1}{n_{\ell}}
$$

where $L_{0}=L-2 M, L_{i}=L_{0}+\sum_{j=1}^{i} 2 n_{j}\left(P_{j}-P_{i+1}\right)$ and $P_{\ell+1}=0$.
Since the $N$-point function $\left\langle s_{1}, s_{1}+d_{1}, \ldots, s_{1}+d_{N-1}\right\rangle_{Y}$ does not depend on the specific site $s_{1}$ (because of translational symmetry), we denote

$$
C_{Y}\left(d_{1}, d_{2}, \ldots, d_{N-1}\right) \equiv\left\langle s_{1}, s_{1}+d_{1}, \ldots, s_{1}+d_{N-1}\right\rangle_{Y}
$$

where $1 \leqslant d_{1}<d_{2}<\cdots<d_{N-1}<L$. Note that $C_{Y}(\emptyset)$ denotes the 1-point function $\left\langle s_{1}\right\rangle_{Y}$.

## Proposition 3.

$$
C_{Y}(\emptyset)=\frac{M}{L} .
$$

Proof. Since $\sum_{n=1}^{L} f(n)=M$,

$$
L C_{Y}(\emptyset)=\sum_{s_{1}=1}^{L}\left\langle s_{1}\right\rangle_{Y}=\frac{1}{\left|\Omega_{Y}\right|} \sum_{f \in \Omega_{Y}} \sum_{n=1}^{L} f(n)=\frac{1}{\left|\Omega_{Y}\right|}\left|\Omega_{Y}\right| M=M .
$$

Next, we consider the 2-point functions.

## Proposition 4.

$$
C_{Y}(1)=\frac{M-s}{L}
$$

where $s=\sum_{i=1}^{\ell} n_{i}$.
Proof. Since $\sum_{n=1}^{L} f(n) f(n+1)=M-s$,

$$
L C_{Y}(1)=\frac{1}{\left|\Omega_{Y}\right|} \sum_{f \in \Omega_{Y}} \sum_{n=1}^{L} f(n) f(n+1)=M-s
$$

In order to investigate $C_{Y}(2)$, let us put

$$
\begin{aligned}
& k_{i}:= \begin{cases}n_{j} & \left(i=P_{j}\right), \\
0 & \text { otherwise },\end{cases} \\
& \hat{k}_{i}:=\sum_{j=i}^{P_{1}} k_{j}, \\
& \tilde{L}:=L-2 \hat{k}_{1} \quad(=L-2 s), \\
& N_{Y}(2):=\sum_{i=3}^{P_{1}} k_{i}(i-2) .
\end{aligned}
$$

Note that $k_{i}$ denotes the number of the rows of the conserved quantity $Y$ with length $i$, and $\hat{k}_{i}$ denotes the number of the rows of the conserved quantity $Y$ with length $\geqslant i$. Thus, $\hat{k}_{1}$ is the number of boxes in the first row of $Y$.

We also define

$$
\begin{aligned}
& V_{f_{0}}:=\left\{f \in \Omega_{Y} \mid E f=f_{0}\right\}, \\
& G_{2}(f):=\sharp\{n \in[L] \mid f(n) f(n+2)=1\} .
\end{aligned}
$$

The following lemma is the key to evaluating $C_{Y}(2)$.
Lemma 1. Let

$$
V_{f_{0}}^{(j)}:=\left\{f \in V_{f_{0}} \mid G_{2}(f)=N_{Y}(2)+j\right\} .
$$

Then, if $V_{f_{0}} \neq \phi, V_{f_{0}}=\bigsqcup_{j=0}^{k_{1}} V_{f_{0}}^{(j)}$ and

$$
\begin{equation*}
\left|V_{f_{0}}^{\left(k_{1}-j\right)}\right|=\frac{v_{j}}{k_{1}!} \tag{7}
\end{equation*}
$$

where

$$
\begin{aligned}
v_{j}:= & \left(\prod_{i=0}^{j-1}\left(\tilde{L}-2 \hat{k}_{2}-i\right)\right)\left(\prod_{i=0}^{k_{1}+j-1}\left(2 \hat{k}_{2}+i\right)\right) \\
& \times\left(\sum_{\substack{0 \leq i_{1}<\ldots<i_{j}<k_{1}+j-1 \\
i_{r}+1 i_{r+1}(r \in[j-1))}} \prod_{h=1}^{j} \frac{1}{\left(2 \hat{k}_{2}+i_{h}\right)\left(2 \hat{k}_{2}+i_{h}+1\right)}\right) .
\end{aligned}
$$

Proof. When $f \in V_{f_{0}}$, there exists a set of positive numbers $\left\{j_{i}\right\}_{i=1}^{k_{1}}\left(1<j_{1}<j_{2}<\cdots<\right.$ $\left.j_{k_{1}} \leqslant \tilde{L}+k_{1}\right)$ such that

$$
f=I\left(j_{1}, j_{2}, \ldots, j_{k_{1}}\right) f_{0}
$$

By examining the positions of 101 and 111, we find that

$$
G_{2}(f)=N_{Y}(2)+\gamma+J
$$

where $\gamma=\gamma\left(f_{0} ;\left\{j_{i}\right\}_{i=1}^{k_{1}}\right)$ is the number of 10 s inserted into the positions adjacent to consecutive 1 s , and $J=\sharp\left\{i \in[d-1] \mid j_{i}+1=j_{i+1}\right\}$ (see figure 2). For example,

$$
f_{0}=001110000100110000
$$

| $f_{0}$ | 00111000 |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $G_{2}\left(f_{0}\right)$ | 1 |  |  |  |
| $f=I(k) f_{0}$ | $\begin{gathered} 001011110000 \\ (k=7) \end{gathered}$ | $\begin{gathered} 001 \underline{1011} \begin{array}{c} 10000 \\ (k=6) \end{array} \end{gathered}$ | $\begin{gathered} 001111010000 \\ (k=4) \end{gathered}$ | $\begin{gathered} 00111000100 \\ (k=2) \end{gathered}$ |
| $G_{2}(f)$ | $\begin{gathered} 3 \\ (\gamma=1, J=0) \end{gathered}$ | $\begin{gathered} 2 \\ (\gamma=0, J=0) \end{gathered}$ | $\begin{gathered} 3 \\ (\gamma=1, J=0) \end{gathered}$ | $\begin{gathered} 2 \\ (\gamma=0, J=0) \end{gathered}$ |

Figure 2. Example $G_{2}(f), \gamma$ and $J$.
and $f=I(5,6,14,15,18) f_{0}$; then

$$
\begin{aligned}
& f=0011110100010100011000111010100000 \\
& (=00111 \underline{10} \underline{100010} \underline{10001} 1000111010100000) .
\end{aligned}
$$

In this example, $k_{1}=5, \hat{k}_{2}=3, N_{Y}(2)=3, \gamma=2$ and $J=2$. Since $0 \leqslant \gamma+J \leqslant k_{1}$, we have the decomposition $V_{f_{0}}=\bigsqcup_{j=0}^{k_{1}} V_{f_{0}}^{(j)}$.

To know $\left|V_{f_{0}}^{(j)}\right|$, we only have to count the number of states with $\gamma+J=j$.
For $k_{1}=1,\left|V_{f_{0}}\right|=\tilde{L}$. Since there are $\hat{k}_{2}$ sets of consecutive $1 \mathrm{~s}, 2 \hat{k}_{2}$ states have $\gamma+J=1$ $(\gamma=1, J=0)$ and the other $\tilde{L}-2 \hat{k}_{2}$ states have $\gamma+J=0(\gamma=0, J=0)$.

For $k_{1}=2$, as was seen in case $k_{1}=1$, there are $2 \hat{k}_{2}$ positions at which $\gamma+J$ can be increased by 1. If one 10 pair is inserted in one of these positions, then there are $2 \hat{k}_{2}+1$ positions for the other pair to increase $\gamma+J$ by 1 and $\tilde{L}-2 \hat{k}_{2}$ positions not to increase it. On the other hand, if one 10 pair is inserted at one of the $\tilde{L}-2 \hat{k}_{2}$ non-increasing positions, then there are $2 \hat{k}_{2}+2$ positions for the other pair to increase $\gamma+J$ by 1 and $\tilde{L}-2 \hat{k}_{2}-1$ positions not to increase it. Hence, considering duplication of insertion, there are $\left(2 \hat{k}_{2}\right)\left(2 \hat{k}_{2}+1\right) / 2$ ! states with $\gamma+J=2$, $\left[\left(2 \hat{k}_{2}\right)\left(\tilde{L}-2 \hat{k}_{2}\right)+\left(\tilde{L}-2 \hat{k}_{2}\right)\left(2 \hat{k}_{2}+2\right)\right] / 2$ ! states with $\gamma+J=1$ and $\left(\tilde{L}-2 \hat{k}_{2}\right)\left(\tilde{L}-2 \hat{k}_{2}-1\right) / 2$ ! states with $\gamma+J=0$.

In general, we can proceed in a similar manner and, referring to the chart in figure 3 , we obtain (7).

## Proposition 5.

$$
C_{Y}(2)=\frac{\sum_{j=0}^{k_{1}} v_{j}\left(\sum_{i=3}^{P_{1}} k_{i}(i-2)+\left(k_{1}-j\right)\right)}{L \sum_{j=0}^{k_{1}} v_{j}}
$$

Proof. From lemma 1, we see that if $V_{f_{0}} \neq \phi$,

$$
\sum_{f \in V_{f_{0}}} \sum_{n=1}^{L} f(n) f(n+2)=\sum_{j=0}^{k_{1}} \frac{v_{j}}{k_{1}!}\left(N_{Y}(2)+\left(k_{1}-j\right)\right)
$$

and

$$
\left|V_{f_{0}}\right|=\sum_{j=0}^{k_{1}} \frac{v_{j}}{k_{1}!}
$$



Figure 3. A chart corresponding to $\gamma+J$ in the proof of lemma 1.

Since the right-hand side of the last equation does not depend on $f_{0}$, and since any state $f \in \Omega_{Y}$ belongs to some $V_{f_{0}}$, we obtain

$$
L C_{Y}(2)=\frac{1}{\left|\Omega_{Y}\right|} \sum_{f \in \Omega_{Y}} \sum_{n=1}^{L} f(n) f(n+2)=\frac{\sum_{j=0}^{k_{1}} v_{j}\left(N_{Y}(2)+\left(k_{1}-j\right)\right)}{\sum_{j=0}^{k_{1}} v_{j}}
$$

For $C_{Y}(d)(d \geqslant 3)$, we can use similar arguments based on elementary combinatorics. However, the expressions become more and more complicated when the difference $d$ increases. Instead, in the next section, we shall use proposition 1 to obtain expressions for general $N$-point functions.

## 3. $N$-point correlation functions for the PBBS

Let the state $f_{0}$ and the set $\mathcal{X}_{Y} \subset \mathbb{Z}_{+}^{n_{1}} \times \mathbb{Z}_{+}^{n_{2}} \times \cdots \times \mathbb{Z}_{+}^{n_{\ell}}\left(=\mathbb{Z}_{+}^{s}\right)$ be

$$
f_{0}=\underbrace{000 \cdots 00}_{L_{0}}
$$

and

$$
\mathcal{X}_{Y}:=\left\{\begin{array}{r|r}
\left\{x_{i}(k)\right\}_{i=1, k=1}^{\ell,} n_{i} & 1<x_{i}(1)<x_{i}(2)<\cdots<x_{i}\left(n_{i}\right) \leqslant L_{i-1}+n_{i}  \tag{8}\\
(i=1,2, \ldots, \ell)
\end{array}\right\} .
$$

We define the state $f_{X}$ recursively as

$$
\begin{aligned}
f_{j} & :=\underbrace{I(\emptyset) \cdots I(\emptyset)}_{P_{j}-P_{j+1}-1} I\left(X_{j}\right) f_{j-1} \quad(j=1,2, \ldots, \ell), \\
f_{X} & :=f_{\ell}
\end{aligned}
$$

where $X_{j}=\left\{x_{j}(k)\right\}_{k=1}^{n_{j}} \subset X \in \mathcal{X}_{Y}$. Note that, from the definition of a 10-insertion, $I(\emptyset)$ is the procedure needed to insert 10 s between 10:

$$
\begin{aligned}
& f=00111001111000001101000111000000, \\
& I(\emptyset) f=00111 \boxed{10} 00111 \boxed{10} 0000011 \boxed{10} 01 \boxed{10} 000111 \boxed{10} 000000,
\end{aligned}
$$

and $f_{X} \in \Omega_{Y}$ by construction. We also define $\tilde{\Omega}_{Y}$ by

$$
\tilde{\Omega}_{Y}:=\left\{f_{X} \mid X \in \mathcal{X}_{Y}\right\} .
$$

## Lemma 2.

$$
\begin{equation*}
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle_{Y}=\frac{1}{L\left|\tilde{\Omega}_{Y}\right|} \sum_{f \in \tilde{\Omega}_{Y}} \sum_{k=1}^{L} f\left(k+s_{1}\right) f\left(k+s_{2}\right) \cdots f\left(k+s_{N}\right) . \tag{9}
\end{equation*}
$$

Proof. By virtue of the definition of $f_{X}, \tilde{\Omega}_{Y}$ is the set of states with conserved quantities $Y$ and the last entry of the 10 -sequence is one of the 0 s that are not marked in the time evolution rule, i.e. $f_{X}(L)=\left(T_{L} f_{X}\right)(L)=0$. By defining the shift operator $S$ by $(S f)(n):=f(n+1)$, and $\left(S^{k} f\right):=S\left(S^{k-1} f\right)(k=1,2, \ldots)$ with $S^{0} f:=f$ and for sets

$$
S^{k} \tilde{\Omega}_{Y}:=\left\{S^{k} f_{X} \mid X \in \mathcal{X}_{Y}\right\} \quad(k=1,2, \ldots, L)
$$

we find

$$
\forall f \in \Omega_{Y}, \quad \sharp\left\{k \mid f \in S^{k} \tilde{\Omega}_{Y}(k=1,2, \ldots, L)\right\}=L_{0} .
$$

Note that $S^{L} f=f$. Since $\left|\Omega_{Y}\right|=\frac{L}{L_{0}}\left|\tilde{\Omega}_{Y}\right|$,

$$
\begin{aligned}
\left\langle s_{1}, s_{2}, \ldots, s_{N}\right\rangle_{Y} & =\frac{1}{\left|\Omega_{Y}\right|} \sum_{f \in \Omega_{Y}} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right) \\
& =\frac{1}{L\left|\tilde{\Omega}_{Y}\right|} \sum_{k=1}^{L} \sum_{f \in S^{k} \tilde{\Omega}_{Y}} f\left(s_{1}\right) f\left(s_{2}\right) \cdots f\left(s_{N}\right) \\
& =\frac{1}{L\left|\tilde{\Omega}_{Y}\right|} \sum_{k=1}^{L} \sum_{f \in \tilde{\Omega}_{Y}} f\left(s_{1}+k\right) f\left(s_{2}+k\right) \cdots f\left(s_{N}+k\right) .
\end{aligned}
$$

Thus, we obtain (9).
Proposition 6. For $X \in \mathcal{X}_{Y}, f_{X}$ is explicitly given as

$$
f_{X}(n)=u_{n}^{0}(X)
$$

where

$$
\begin{align*}
u_{n}^{t}(X):= & \eta_{n+1}^{t-1}(X)-\eta_{n+1}^{t}(X)-\eta_{n}^{t-1}(X)+\eta_{n}^{t}(X), \\
\eta_{n}^{t}(X):= & \max _{\substack{m_{j i j} \in \mathbb{Z}, i \in\left[l_{j} ; \in \in n_{i}\right]}}\left[\sum_{i=1}^{\ell} \sum_{k=1}^{n_{i}} m_{i k}\left(t P_{i}-n-x_{i}(k)+L+k+1+\frac{Z_{i}}{2}\right)\right. \\
& \left.-\sum_{i=1}^{\ell} \sum_{k=1}^{n_{i}} \sum_{j=1}^{\ell} \sum_{h=1}^{n_{j}} m_{i k} \Xi_{i k j h} m_{j h}\right], \tag{10}
\end{align*}
$$

$$
\begin{aligned}
& \Xi_{i k j h}:=\frac{Z_{i}}{2} \delta_{i j} \delta_{k h}+P_{\max [i, j]}, \\
& Z_{i}:=L-2\left(P_{i} \sum_{j=1}^{i} n_{j}+\sum_{j=i+1}^{\ell} n_{j} P_{j}\right) .
\end{aligned}
$$

Proof. From proposition $1, f_{X}$ is determined by the parameters $W_{n}$ and $a_{n}(n=1,2, \ldots, s)$. Here $W_{n}$ is the amplitude of the $n$th soliton and $a_{n}$ is its position, i.e. the position of the $n$th 10 -wall, counting from the right. From the definition of the position and the amplitude of a soliton, it follows that both can be determined from 10 -insertions. Because of the way $f_{X}$ was constructed, the set $\left\{x_{j}(k)\right\}_{k=1}^{n_{j}}$ corresponds to the position of $n_{j}$ solitons with amplitude $P_{j}$, though it does not directly give their position. Hereafter, we shall refer to a soliton with amplitude $P$ as a $P$-soliton. By considering the relation between the position of a soliton and 10 -insertions, we find that the position of the $k$ th $P_{j}$-soliton counting from the right is $L-x_{j}^{(\ell)}(k)+2$, where $x_{j}^{(\ell)}(k)$ is determined recursively: we define $x_{j}^{(i)}(k)(i \in[\ell], j \in$ [ $i], k \in\left[n_{j}\right]$ ) as

$$
\begin{aligned}
x_{j}^{(i)}(k):= & x_{j}(k)+\left(P_{j}-P_{i+1}\right)\left(2 \beta_{j}(k)+2 k-1\right) \\
& +\sum_{s=j+1}^{i} 2\left(P_{s}-P_{i+1}\right) \alpha_{j}^{(s)}(k)-k+1,
\end{aligned}
$$

where
$\alpha_{j}^{(i)}(k):=\sharp\left\{r \in\left[n_{i}\right] \mid L_{i-1}+n_{i}-x_{i}(r)+1>g_{j}^{(i)}(k)\right\}$,
$\beta_{1}(k):=0, \quad \beta_{i}(k):=\sum_{s=1}^{i-1} \sharp\left\{r \in\left[n_{s}\right] \mid g_{s}^{(i)}(r) \geqslant L_{i-1}+n_{i}-x_{i}(k)+1\right\}$,
$g_{j}^{(i)}(k):=\max \left\{m \in\left[L_{i-1}+n_{i}\right] \left\lvert\, \begin{array}{rl}m= & L_{i-1}-x_{j}^{(i-1)}(k)+1 \\ & +\sharp\left\{r \in\left[n_{i}\right] \mid L_{i-1}+n_{i}-x_{i}(r)+1 \leqslant m\right\}\end{array}\right.\right\}$.
Note that $x_{j}^{(i)}(1)<x_{j}^{(i)}(2)<\cdots<x_{j}^{(i)}\left(n_{j}\right)$. We show an example in appendix A.
Recalling the fact that $\forall\left\{r \in[d] \mid L+d-j_{r}+1<g_{k}^{\prime}\right\}$ in (4) is the number of inserted $\underline{10}$ s, on the left of the $k$ th soliton (here we do not count the inserted $\underline{10}$ s as solitons), the concrete meaning of these variables becomes clear: $\alpha_{j}^{(i)}(k)$ denotes the number of $P_{i}$-solitons on the right of the $k$ th $P_{j}$-soliton, and $\beta_{j}(k)$ denotes the number of solitons with amplitudes greater than $P_{j}$, to the right of the $k$ th $P_{j}$-soliton.

Since $\left\{L-x_{j}^{(\ell)}(k)+2\right\}_{j=1, k=1}^{\ell, n_{j}}$ is the complete set of positions of the solitons, there exists a one-to-one mapping $\rho:\left\{(j, k) \mid j \in[\ell], k \in\left[n_{j}\right]\right\} \rightarrow[s]$ such that

$$
a_{\rho(j, k)}=L-x_{j}^{(\ell)}(k)+2 .
$$

From these recursion relations, we have

$$
\begin{aligned}
x_{j}^{(\ell)}(k) & =x_{j}(k)+P_{j}\left(2 \beta_{j}(k)+2 k-1\right)+\sum_{i=j+1}^{\ell} 2 P_{i} \alpha_{j}^{(i)}(k)-k+1 \\
& =x_{j}(k)+2\left\{P_{j}\left(\beta_{j}(k)+(k-1)\right)+\sum_{i=j+1}^{\ell} P_{i} \alpha_{j}^{(i)}(k)\right\}+P_{j}-k+1
\end{aligned}
$$

Since the position of the $k$ th $P_{j}$-soliton is $a_{\rho(j, k)}, W_{\rho(j, k)}=P_{j}$ and the set of amplitudes of the solitons on the right of the $k$ th $P_{j}$-soliton is nothing but $\left\{W_{h}\right\}_{h=1}^{\rho(j, k)-1}$. From the definition of $\alpha_{j}^{(i)}(k), \beta_{j}(k)$,

$$
\begin{aligned}
& \alpha_{j}^{(i)}(k)=\sharp\left\{h \in[\rho(j, k)-1] \mid W_{h}=P_{i}\right\}, \\
& \beta_{j}(k)=\sharp\left\{h \in[\rho(j, k)-1] \mid W_{h}>P_{j}\right\},
\end{aligned}
$$

and

$$
\sharp\left\{h \in[\rho(j, k)-1] \mid W_{h}=P_{j}\right\}=k-1 .
$$

Thus, we obtain

$$
x_{j}^{(\ell)}(k)=x_{j}(k)+\sum_{h=1}^{\rho(j, k)-1} 2 \min \left\{W_{\rho(j, k)}, W_{h}\right\}+W_{\rho(j, k)}-k+1 .
$$

Therefore, we find a concrete expression of $a_{\rho(j, k)}$, and (10) is immediately obtained from (2) and (3).

From lemma 2 and proposition 6, we immediately obtain the following theorem.
Theorem 1. Let $\mathcal{X}_{Y}$ be the set defined in (8). We then have

$$
C_{Y}\left(d_{1}, d_{2}, \ldots, d_{N-1}\right)=\frac{1}{L\left|\mathcal{X}_{Y}\right|} \sum_{X \in \mathcal{X}_{Y}} \sum_{n=1}^{L} u_{n}(X) \prod_{i=1}^{N-1} u_{n+d_{i}}(X),
$$

for $u_{n}(X) \equiv u_{n}^{0}(X)$, as given in (10).

## 4. Concluding remarks

In this paper, we investigated correlation functions for the PBBS and obtained explicit forms for 1-point and 2-point functions at short distances. We also give expressions in terms of ultradiscrete theta functions for general $N$-point functions. Investigating their asymptotic properties and clarifying the relation to correlation functions for quantum integrable systems are problems that will be addressed in the future.

Finally, we comment on the time averages of quantities in the PBBS. The time average

$$
C_{f}\left(d_{1}, d_{2}, \ldots, d_{N-1}\right)=\frac{1}{L\left|\mathcal{T}_{f}\right|} \sum_{t=1}^{\mathcal{T}_{f}} \sum_{n=1}^{L}\left(T_{L}^{t} f\right)(n) \prod_{j=1}^{N-1}\left(T_{L}^{t} f\right)\left(n+d_{j}\right)
$$

where $\mathcal{T}_{f}$, the fundamental cycle of $f \in \Omega_{L}$, depends not only on the conserved quantities of the state but, in general, also on the initial state $f$ itself. For example, the conserved quantities of the states $f_{1}=0100100$ and $f_{2}=0101000$ are the same, but $C_{f_{1}}(3)=\frac{1}{7}$ and $C_{f_{2}}(3)=0$. Hence, in general, $C_{f}\left(d_{1}, d_{2}, \ldots, d_{N-1}\right) \neq C_{Y}\left(d_{1}, d_{2}, \ldots, d_{N-1}\right)$ even for $f \in \Omega_{Y}$. Note that for the 1-point function $C_{f}(\emptyset)$, we can easily show that

$$
\forall f \in \Omega_{Y}, \quad C_{f}(\emptyset)=C_{Y}(\emptyset)=\frac{M}{L}
$$

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## Appendix A. Example to determine the position of solitons

Let $L=33,\left(P_{1}, n_{1}\right)=(3,2),\left(P_{2}, n_{2}\right)=(2,3)$ and $\left(P_{3}, n_{3}\right)=(1,2) . \quad$ Since $L_{0}=5, L_{1}=9, L_{2}=19$, we suppose $\left\{x_{i}(k)\right\}_{i=1, k=1}^{\ell, n_{i}} \in \mathcal{X}_{Y}$ to be
$x_{1}(1)=3, x_{1}(2)=6 ; x_{2}(1)=2, x_{2}(2)=5, x_{2}(3)=9 ; x_{3}(1)=12, x_{3}(2)=17$.
In this case, by successive 10 -insertions, we obtain the state

$$
011100100110001001110011000011000 .
$$

Then, we see from [9] that

$$
\left\{\left(a_{k}, W_{k}\right)\right\}_{k=1}^{7}=\{(31,2),(25,3),(21,2),(16,1),(12,2),(8,1),(5,3)\}
$$

On the other hand, according to the algorithm in the proof of proposition 6, the variables are determined as follows.
(i) The variables $x_{1}^{(1)}(k)$ are determined from the initial data:

$$
\begin{aligned}
x_{1}^{(1)}(1) & =x_{1}(1)+\left(P_{1}-P_{2}\right)\left(2 \beta_{1}(1)+2 \cdot 1-1\right)-1+1 \\
& =3+(3-2)(2 \cdot 0+2 \cdot 1-1)-1+1=4, \\
x_{1}^{(1)}(2) & =x_{1}(2)+\left(P_{1}-P_{2}\right)\left(2 \beta_{1}(2)+2 \cdot 2-1\right)-2+1 \\
& =6+(3-2)(2 \cdot 0+2 \cdot 2-1)-2+1=8 .
\end{aligned}
$$

(ii) To determine the variables $x_{j}^{(2)}(k)$, we calculate $g_{1}^{(2)}(k), \alpha_{j}^{(2)}(k)$ and $\beta_{2}(k)$ :

$$
\begin{aligned}
& g_{1}^{(2)}(1)=\max \left\{m \in\left[L_{1}+n_{2}\right] \left\lvert\, \begin{array}{rl}
m= & L_{1}-x_{1}^{(1)}(1)+1 \\
& +\sharp\left\{r \in\left[n_{2}\right] \mid L_{1}+n_{2}-x_{2}(r)+1 \leqslant m\right\}
\end{array}\right.\right\} \\
& \left.=\max \left\{m \in[12] \left\lvert\, m=6+\sharp\left\{r \in[3] \left\lvert\, \begin{array}{cc}
11 & (r=1) \\
8 & (r=2) \\
4 & (r=3)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=8 \text {, } \\
& \left.g_{1}^{(2)}(2)=\max \left\{m \in[12] \left\lvert\, m=2+\sharp\left\{r \in[3] \left\lvert\, \begin{array}{cc}
11 & (r=1) \\
8 & (r=2) \\
4 & (r=3)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=2 \text {; } \\
& \alpha_{1}^{(2)}(1)=\sharp\left\{r \in\left[n_{2}\right] \mid L_{1}+n_{2}-x_{2}(r)+1>g_{1}^{(2)}(1)\right\} \\
& \left.=\sharp\left\{r \in[3] \left\lvert\, \begin{array}{cc}
11 & (r=1) \\
8 & (r=2) \\
4 & (r=3)
\end{array}\right.\right\}>8\right\}=1, \\
& \left.\alpha_{1}^{(2)}(2)=\sharp\left\{r \in[3] \left\lvert\, \begin{array}{cc}
11 & (r=1) \\
8 & (r=2) \\
4 & (r=3)
\end{array}\right.\right\}>2\right\}=3 \text {; } \\
& \beta_{2}(1)=\sharp\left\{r \in\left[n_{1}\right] \mid g_{1}^{(2)}(r)>L_{1}+n_{2}-x_{2}(1)+1\right\} \\
& \left.=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
8 & (r=1) \\
2 & (r=2)
\end{array}\right.\right\} \geqslant 11\right\}=0,
\end{aligned}
$$

$$
\begin{aligned}
& \left.\beta_{2}(2)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
8 & (r=1) \\
2 & (r=2)
\end{array}\right.\right\} \geqslant 8\right\}=1, \\
& \left.\beta_{2}(3)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
8 & (r=1) \\
2 & (r=2)
\end{array}\right.\right\} \geqslant 4\right\}=1 .
\end{aligned}
$$

Then, we have

$$
\begin{aligned}
x_{1}^{(2)}(1) & =x_{1}(1)+\left(P_{1}-P_{3}\right)\left(2 \beta_{1}(1)+2 \cdot 1-1\right)+2\left(P_{2}-P_{3}\right) \alpha_{1}^{(2)}(1)-1+1 \\
& =3+(3-1)(2 \cdot 0+2 \cdot 1-1)+2(2-1) \cdot 1-1+1=7, \\
x_{1}^{(2)}(2) & =17 \\
x_{2}^{(2)}(1) & =x_{2}(1)+\left(P_{2}-P_{3}\right)\left(2 \beta_{2}(1)+2 \cdot 1-1\right)-1+1 \\
& =2+(2-1)(2 \cdot 0+2 \cdot 1-1)-1+1=3, \\
x_{2}^{(2)}(2) & =9, \quad x_{2}^{(2)}(3)=14 .
\end{aligned}
$$

(iii) To determine the variables $x_{j}^{(3)}(k)$, we calculate $g_{j}^{(3)}(k), \alpha_{j}^{(3)}(k)$ and $\beta_{3}(k)$ :

$$
\begin{aligned}
& \left.g_{1}^{(3)}(1)=\max \left\{m \in[21] \left\lvert\, m=13+\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=15, \\
& \left.g_{1}^{(3)}(2)=\max \left\{m \in[21] \left\lvert\, m=3+\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=3 ; \\
& \left.g_{2}^{(3)}(1)=\max \left\{m \in[21] \left\lvert\, m=17+\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=19, \\
& \left.g_{2}^{(3)}(2)=\max \left\{m \in[21] \left\lvert\, m=11+\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=13, \\
& \left.g_{2}^{(3)}(3)=\max \left\{m \in[21] \left\lvert\, m=6+\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\} \leqslant m\right.\right\}\right\}=7 ; \\
& \left.\alpha_{1}^{(3)}(1)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\}>15\right\}=0, \\
& \left.\alpha_{1}^{(3)}(2)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\}>3\right\}=2 ; \\
& \left.\alpha_{2}^{(3)}(1)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\}>19\right\}=0, \\
& \left.\alpha_{2}^{(3)}(2)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\}>13\right\}=0, \\
& \left.\alpha_{2}^{(3)}(3)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{ll}
10 & (r=1) \\
5 & (r=2)
\end{array}\right.\right\}>7\right\}=1 ;
\end{aligned}
$$

$$
\begin{aligned}
& \left.\left.\beta_{3}(1)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
15 & (r=1) \\
3 & (r=2)
\end{array}\right.\right\} \geqslant 10\right\}+\sharp\left\{r \in[3] \left\lvert\, \begin{array}{cc}
19 & (r=1) \\
13 & (r=2) \\
7 & (r=3)
\end{array}\right.\right\} \geqslant 10\right\}=3, \\
& \left.\left.\beta_{3}(2)=\sharp\left\{r \in[2] \left\lvert\, \begin{array}{cc}
15 & (r=1) \\
3 & (r=2)
\end{array}\right.\right\} \geqslant 5\right\}+\sharp\left\{r \in[3] \left\lvert\, \begin{array}{cc}
19 & (r=1) \\
13 & (r=2) \\
7 & (r=3)
\end{array}\right.\right\} \geqslant 5\right\}=4 .
\end{aligned}
$$

Then, we have

$$
\begin{aligned}
x_{1}^{(3)}(1) & =x_{1}(1)+P_{1}\left(2 \beta_{1}(1)+2 \cdot 1-1\right)+\sum_{s=2}^{3} 2 P_{s} \alpha_{1}^{(s)}(1)-1+1 \\
& =3+3(2 \cdot 0+2 \cdot 1-1)+2(2 \cdot 1+1 \cdot 0)-1+1=10, \\
x_{1}^{(3)}(2) & =30 ; \\
x_{2}^{(3)}(1) & =x_{2}(1)+P_{2}\left(2 \beta_{2}(1)+2 \cdot 1-1\right)+2 P_{3} \alpha_{2}^{(3)}(1)-1+1 \\
& =2+2(2 \cdot 0+2 \cdot 1-1)+2 \cdot 1 \cdot 0-1+1=4, \\
x_{2}^{(3)}(2) & =14, \quad x_{2}^{(3)}(3)=23 ; \\
x_{3}^{(3)}(1) & =x_{3}(1)+P_{3}\left(2 \beta_{3}(1)+2 \cdot 1-1\right)-1+1 \\
& =12+1 \cdot(2 \cdot 3+2 \cdot 1-1)-1+1=19, \\
x_{3}^{(3)}(2) & =27 .
\end{aligned}
$$

(iv) Finally, we obtain the position of solitons:

$$
L-x_{j}^{(3)}(k)+2=\left\{\begin{array}{cl}
25 & (j=1, k=1) \\
5 & (j=1, k=2) \\
31 & (j=2, k=1) \\
21 & (j=2, k=2) \\
12 & (j=2, k=3) \\
16 & (j=3, k=1) \\
8 & (j=3, k=2)
\end{array}\right.
$$

Therefore, we confirm that $L-x_{j}^{(\ell)}(k)+2$ is the position of the $k$ th soliton among the solitons with amplitude $P_{j}$.

## Appendix B. Example of values for the correlation function

From theorem 1, we obtain the following examples.
(a) $L=12 ; P_{1}=3, n_{1}=1 ; P_{2}=1, n_{2}=2:$
$C_{Y}(\emptyset)=\frac{5}{12}, \quad C_{Y}(1)=\frac{1}{6}, \quad C_{Y}(2)=\frac{13}{84}, \quad C_{Y}(3)=\frac{19}{126}, \quad C_{Y}(1,2)=\frac{5}{84} ;$
(b) $L=14 ; P_{1}=2, n_{1}=2 ; P_{2}=1, n_{2}=2$ :
$C_{Y}(\emptyset)=\frac{3}{7}, \quad C_{Y}(1)=\frac{1}{7}, \quad C_{Y}(2)=\frac{5}{49}, \quad C_{Y}(3)=\frac{82}{441}, \quad C_{Y}(1,2)=0 ;$
(c) $L=14 ; P_{1}=3, n_{1}=1 ; P_{2}=1, n_{2}=3$ :
$C_{Y}(\emptyset)=\frac{3}{7}, \quad C_{Y}(1)=\frac{1}{7}, \quad C_{Y}(2)=\frac{5}{28}, \quad C_{Y}(3)=\frac{69}{392}, \quad C_{Y}(1,2)=\frac{5}{112} ;$
(d) $L=14 ; P_{1}=3, n_{1}=1 ; P_{2}=2, n_{2}=1 ; P_{3}=1, n_{3}=1$ :
$C_{Y}(\emptyset)=\frac{3}{7}$,
$C_{Y}(1)=\frac{3}{14}$,
$C_{Y}(2)=\frac{3}{28}$,
$C_{Y}(3)=\frac{13}{112}$,
$C_{Y}(1,2)=\frac{1}{16}$.

## References

[1] Baxter R J 1982 Exactly Solved Models in Statistical Mechanics (London: Academic)
[2] For example, Boos H, Jimbo M, Miwa T, Smirnov F and Takeyama Y 2006 Algebraic representation of correlation functions in integrable spin chains Ann. Henri Poincaré 7 1395-428
[3] Takahashi D and Satsuma J 1990 A soliton cellular automaton J. Phys. Soc. Japan 59 3514-9
[4] Yura F and Tokihiro T 2002 On a periodic soliton cellular automaton J. Phys. A: Math. Gen. 35 3787-801
[5] Fukuda K, Okado M and Yamada Y 2000 Energy functions in box ball systems Int. J. Mod. Phys. A 15 1379-92
[6] Hatayama G, Hikami K, Inoue R, Kuniba A, Takagi T and Tokihiro T 2001 The $A_{M}^{(1)}$ automata related to crystals of symmetric tensors J. Math. Phys. 42 274-308
[7] Mada J, Idzumi M and Tokihiro T 2006 The exact correspondence between conserved quantities of a periodic box-ball system and string solutions of the Bethe ansatz equations J. Math. Phys. 47053507
[8] Torii M, Takahashi D and Satsuma J 1996 Combinatorial representation of invariants of a soliton cellular automaton Physica D 92 209-20
[9] Mada J, Idzumi M and Tokihiro T 2008 The box-ball system and the $N$-soliton solution of the ultradiscrete KdV equation J. Phys. A: Math. Theor. 41175207
[10] Mikhalkin G B and Zharkov I 2008 Tropical curves, their Jacobians and theta functions Curves and Abelian Varieties, Contemp. Math. 465 203-30
[11] Yoshihara D, Yura F and Tokihiro T 2003 Fundamental cycle of a periodic box-ball system J. Phys. A: Math. Gen. 36 99-121

